

Structure and Properties of Triclinic $\text{Ni}_{0.85}\text{Mo}_6\text{Te}_8$ *

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The structure of $\text{Ni}_{0.85}\text{Mo}_6\text{Te}_8$ was refined from single-crystal X-ray diffraction data at room temperature. It is triclinic, space group $P\bar{1}$, $a = 7.028(2)$ Å, $b = 7.100(2)$ Å, $c = 7.102(2)$ Å, $\alpha = 91.23(2)^\circ$, $\beta = 95.73(2)^\circ$, $\gamma = 90.82(2)^\circ$, $V = 352.48$ Å³, $Z = 1$; 1619 reflections, 75 refined parameters, $R = 0.031$. The Mo atoms form distorted octahedral clusters (2.69 Å $\leq d_{\text{intra}}[\text{Mo-Mo}] \leq 2.81$ Å; 3.58 Å $< d_{\text{inter}}[\text{Mo-Mo}]$). The Ni atoms are disordered (site occupancy: 0.423(7); $d[\text{Ni-Ni}] = 2.586(6)$ Å), and interact strongly with one Mo_6 cluster ($d[\text{Ni-Mo}] = 2.603(3)$ and $2.958(3)$ Å), and weakly with another ($d[\text{Ni-Mo}] = 2.985(3)$ Å). The structure transforms at 1057(5) K into a rhombohedral modification ($a_{\text{hex}} = 10.457(2)$ Å, $c_{\text{hex}} = 11.866(3)$ Å at 1073 K). Measurements on powders suggest metallic conductivity (5.1×10^{-4} Ω-cm at 293 K) and weakly temperature-dependent paramagnetism (110×10^{-6} emu/g at 100 K). © 1987 Academic Press, Inc.

1. Introduction

Chevrel phases $M_x\text{Mo}_6X_8$ ($M_x = \text{metal}$, $X = \text{chalcogen}$) based on tellurium are relatively rare. Tellurides have so far only been reported for $M_x = \text{Fe}$ ($0 < x < 2$), Co ($0 < x < 0.66$), Ni ($0 < x < 2$), and Cu ($0 < x < 2$) (1). Their crystal symmetries at room temperature were either rhombohedral (Fe, Co) or triclinic (Ni, Cu). Apart from lattice parameters for the Fe- and Co-containing compounds no other information was given.

In this communication we report on a structural characterization and some properties of triclinic $\text{Ni}_{0.85}\text{Mo}_6\text{Te}_8$. Preliminary results were presented elsewhere (2).

* Dedicated to Dr. H. Nowotny.

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2. Experimental and Results

A sample of nominal composition NiMo_6Te_8 was prepared by reaction of binary Mo_6Te_8 with elemental Ni at 1300 K in sealed quartz ampoules. The preparation of Mo_6Te_8 was described previously (3). X-ray Guinier photographs confirmed the formation of a new ternary compound. The intensity profile was similar to that of rhombohedral Mo_6Te_8 ($a_r = 7.10$ Å, $\alpha_r = 92.60^\circ$; according to Ref. (1); $a_r = 7.050(3)$ Å, $\alpha_r = 92.54(3)^\circ$, $V_r = 349.3$ Å³, according to Ref. (4)), except for line splitting which suggested a triclinic lattice distortion. A single crystal of approximately cubic shape (edge lengths ≈ 0.02 mm) was isolated and examined on an automated four-circle X-ray diffractometer. Cell parameters (see abstract) were determined by least-squares refinement of 12 measured Bragg angles in the region $15^\circ \leq 2\theta < 25^\circ$. Integrated inten-

TABLE I
 ATOMIC PARAMETERS

Atom	<i>x</i>	<i>y</i>	<i>z</i>	<i>U</i> _{eq.}	<i>U</i> ₁₁	<i>U</i> ₃₃	<i>U</i> ₂₂	<i>U</i> ₂₃	<i>U</i> ₁₂	<i>U</i> ₁₃
Mo(1)	0.2423(1)	0.4239(1)	0.5411(1)	76(2)	60(4)	72(4)	95(4)	-5(3)	7(3)	6(3)
Mo(2)	0.4108(1)	0.5499(1)	0.2434(1)	83(2)	65(4)	88(4)	95(4)	-1(3)	1(3)	3(3)
Mo(3)	0.5432(1)	0.2376(1)	0.4228(1)	88(2)	75(4)	99(4)	91(4)	4(3)	5(3)	24(3)
Te(1)	0.1034(1)	0.7463(1)	0.3621(1)	120(2)	86(3)	107(3)	173(4)	24(2)	40(2)	26(2)
Te(2)	0.7344(1)	0.3721(1)	0.1314(1)	105(2)	95(3)	124(3)	95(3)	-8(2)	7(2)	20(2)
Te(3)	0.3973(1)	0.1157(1)	0.7509(1)	142(2)	173(4)	88(3)	176(3)	-26(2)	77(3)	-27(2)
Te(4)	0.2020(1)	0.2237(1)	0.1968(1)	126(2)	94(3)	113(3)	168(3)	-1(2)	-12(2)	-9(2)
Ni	0.1718(4)	0.0645(4)	0.4916(4)	116(9)	87(16)	118(16)	143(16)	-1(11)	6(11)	-1(11)

Note. Ni_{0.85}Mo₆Te₈, *a* = 7.028(2) Å, *b* = 7.100(2) Å, *c* = 7.102(2) Å, α = 91.23(2)°, β = 95.73(2)°, γ = 90.82(2)°, *P*1̄ (No. 2), *Z* = 1, 1619 *hkl*, *d*(calc) = 7.753 g/cm³, *R* = 0.031, all atoms in 2(*a*), *PP*(Ni) = 0.423(7). The *U*_{*ij*} (Å² × 100) are expressed as exp -2π²(*h*²*a*²*U*₁₁ + ··· + 2*hka***b***U*₁₂). *U*_{eq.} is defined as one-third of the trace of the orthogonalized *U*_{*ij*} tensor.

sities of 1619 symmetry-independent reflections were collected at room temperature in the range 3° ≤ 2θ ≤ 55° by using continuous θ - 2θ scans and graphite monochromatized MoKα radiation. Absorption effects (μ[MoKα] = 244 cm⁻¹) were corrected empirically from azimuthal scans of 11 reflections. The structure was solved by direct methods in space group *P*1̄ (380 *E(hkl)* values with *E* ≥ 1.20; 2494 triple-product relations; tests for negative-quartet relations) and refined by conventional least-squares methods. The weighting scheme used was *w* = 1/([σ(*F*)]² + 0.0002*F*²). The position of the Ni atoms was derived from Fourier maps. Preliminary refinements with isotropic temperature factors and fixed occupancy of the Ni atom site (*PP*(Ni) = 0.5) converged at *R* = 0.072. Subsequent refinements with anisotropic temperature factors and variable occupancy of the Ni atom site converged at *R* = 0.031. The refined occupancy factor, *PP*(Ni) = 0.423(7), indicated that the crystal was Ni deficient compared to the overall sample composition and that its stoichiometry was Ni_{0.85(2)}Mo₆Te₈. Refinements based on an ordered Ni atom distribution (*PP*(Ni) = 1.0) in noncentrosymmetric space group

*P*1 gave a poorer fit (*R* = 0.037). All calculations were performed using the program system SHELXTL (5). Atomic parameters are summarized in Table I, and a structural drawing is given in Fig. 1.¹ For comparison with rhombohedral Mo₆Te₈ (space group *R*3̄) the structure was described such that its pseudotrigonal axis runs along [111]. This nonstandardized (6) description corresponds to that adopted in work on other triclinic Chevrel phases (see below).

The stability of Ni_{0.85}Mo₆Te₈ at high temperature was studied by differential thermal analysis (DTA) performed on a NETZSCH DTA 404S apparatus. Approximately 150 mg of finely ground powder was placed in a quartz ampoule which was evacuated and sealed under vacuum. The heating rate was 10 K/min and the maximum heating temperature 1473 K. Upon heating two endothermic peaks were observed, one occurring at 1057(5) K and the other at 1216(5) K. The intensity of the

¹ Additional material to this paper can be ordered from Pachinformationszentrum Energie Physik Mathematik, D-7514 Bggenstein-I.eopoldshafen-2, Federal Republic of Germany, by referring to the authors and title of the paper.

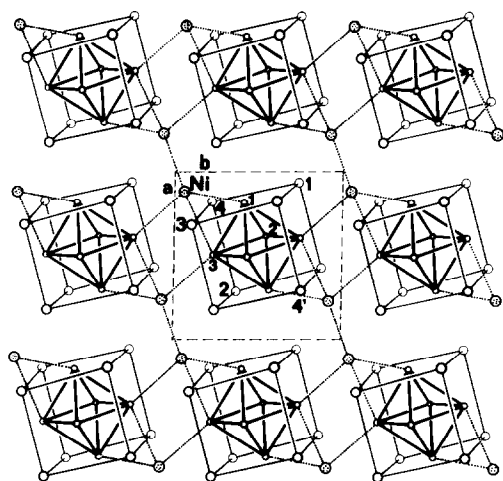


FIG. 1. Projection on a - b plane of $\text{Ni}_{0.85}\text{Mo}_6\text{Te}_8$. Small open circles, Mo; large open circles, Te; dotted circles, Ni atom sites (occupancy: 42%). Atom numbering as in Table I. The pseudo-trigonal axis runs through Te4 and Te4'. Thick lines, Mo-Mo bonds; dotted lines, Mo-Ni and Ni-Ni bonds; thin lines, Te-Te contacts indicating the $[\text{Mo}_6\text{Te}_8]$ units.

latter peak was approximately one-fifth that of the former. High-temperature X-ray powder photographs (Simon-Guinier camera) taken at 1073 and 1243 K suggested that the first DTA peak was due to a structural phase transition leading to a modification of rhombohedral symmetry. Its hexagonal lattice parameters (see abstract) were determined at 1073 K from 19 hkl values in the region $25^\circ \leq 2\theta < 72^\circ$.

The diffraction profile showed good agreement with a theoretical powder pattern calculated for rhombohedral Mo_6Te_8 (4) containing one Ni atom at the inversion center of the chalcogen atom network. The nature of the second DTA peak was not clarified. Upon cooling, both DTA peaks showed a hysteresis of about 13 K, confirming the first-order nature of the structural phase transition. X-ray photographs taken after the DTA experiment showed that the transition was reversible.

The electrical resistivity was measured

on a sintered pellet by the four-point method in the temperature range between 298 and 2.5 K. It decreased continuously from $5.1 \times 10^{-4} \Omega\text{-cm}$ at room temperature to $2.8 \times 10^{-4} \Omega\text{-cm}$ at 2.5 K, suggesting metallic conductivity.

The magnetic susceptibility was measured on a polycrystalline powder sample as a function of temperature ($T = 6$ –300 K, applied field 1000G), and magnetic field (5–50 kG, $T = 5$ –300 K) by using a SQUID magnetometer (7). It decreased from $140 \times 10^{-6} \text{emu/g}$ (300 K) to $110 \times 10^{-6} \text{emu/g}$ (100 K) and then increased to $250 \times 10^{-6} \text{emu/g}$ (4 K). A diamagnetic correction of $-210 \times 10^{-6} \text{emu/g}$ was applied. No variation as a function of applied field was observed. The susceptibility increase below $T = 50$ K was attributed to the presence of about 0.3% of metallic Ni, a value which is consistent with the observed Ni deficiency (0.5 mole%) of the triclinic structure.

3. Discussion

Triclinic $\text{Ni}_{0.85}\text{Mo}_6\text{Te}_8$ is isostructural to its selenium-based congener $\text{Ni}_{0.66}\text{Mo}_6\text{Se}_8$ (8). The structures differ mainly with respect to their triclinic lattice shears relative to their rhombohedral (high-temperature) modification: $a < b$, $b = c$; $\alpha \approx \gamma < \beta$ ($\text{Ni}_{0.85}\text{Mo}_6\text{Te}_8$); $b < a$, $a \approx c$; $\alpha \approx \gamma < \beta$ ($\text{Ni}_{0.66}\text{Mo}_6\text{Se}_8$), thus leading to different standardized (6) structure descriptions (L. Gelato, personal communication). The structure is built up by cube-shaped $[\text{Mo}_6\text{Te}_8]$ units (Fig. 1) which contain triclinic distorted octahedral Mo_6 clusters ($2.688(1) \text{ \AA} < d_{\text{intra}}[\text{Mo-Mo}] \leq 2.812(1) \text{ \AA}$ (telluride); $2.650 \text{ \AA} < d_{\text{intra}}[\text{Mo-Mo}] < 2.818 \text{ \AA}$ (selenide)). The average Mo-Mo bond length within the Mo_6 clusters in the telluride ($d_{\text{av}}[\text{Mo-Mo}] = 2.74 \text{ \AA}$) is almost identical to that in the corresponding Ni-free binary compound Mo_6Te_8 ($d_{\text{av}}[\text{Mo-Mo}] = 2.74 \text{ \AA}$ (4)), whereas that in the selenide ($d_{\text{av}}[\text{Mo-Mo}] = 2.73 \text{ \AA}$) it is

TABLE II
 $\text{Ni}_{0.85}\text{Mo}_6\text{Te}_8$: INTERATOMIC DISTANCES (STANDARD DEVIATION) IN Å

Mo(1)	Mo(2)	Mo(3)	Ni	Te(1)	Te(2)	Te(3)	Te(4)
Mo(2) 2.688(1)	Mo(1) 2.688(1)	Mo(1) 2.705(1)	Mo(1) 2.603(3)	Mo(2) 2.779(1)	Mo(1) 2.703(1)	Mo(2) 2.711(1)	Mo(2) 2.721(1)
Mo(3) 2.705(1)	Mo(3) 2.711(1)	Mo(2) 2.711(1)	Mo(3) 2.958(3)	Mo(3) 2.783(1)	Mo(3) 2.758(1)	Mo(3) 2.787(1)	Mo(3) 2.749(1)
Mo(2) 2.748(1)	Mo(1) 2.748(1)	Mo(2) 2.771(1)	Mo(3) 2.985(3)	Mo(1) 2.784(1)	Mo(2) 2.793(1)	Mo(3) 2.827(1)	Mo(1) 2.790(1)
Mo(3) 2.812(1)	Mo(3) 2.771(1)	Mo(1) 2.812(1)	Ni 2.586(6)	Mo(1) 2.850(1)	Mo(2) 2.823(1)	Mo(1) 2.842(1)	Ni 2.426(3)
Mo(1) 3.584(1)	Mo(2) 3.819(1)	Mo(3) 3.631(1)	Te(3) 2.326(3)	Ni 2.442(3)		Ni 2.326(3)	
Mo(1) 3.869(1)	Mo(2) 3.849(1)	Mo(3) 3.934(1)	Te(4) 2.426(3)	Ni 2.655(3)			
Ni 2.603(3)		Ni 2.958(3)	Te(1) 2.442(3)				
		Ni 2.985(3)	Te(1) 2.655(3)				
Te(2) 2.703(1)	Te(3) 2.711(1)	Te(4) 2.749(1)	Te(3) 3.850(3)				
Te(1) 2.784(1)	Te(4) 2.721(1)	Te(2) 2.758(1)					
Te(4) 2.790(1)	Te(1) 2.779(1)	Te(1) 2.783(1)					
Te(3) 2.842(1)	Te(2) 2.793(1)	Te(3) 2.787(1)					
Te(1) 2.850(1)	Te(2) 2.823(1)	Te(3) 2.827(1)					

smaller by 0.03 Å than that in binary Mo_6Se_8 ($d_{\text{av}}[\text{Mo}-\text{Mo}] = 2.76$ Å (9)) (Table II). The Mo_6 clusters interact only weakly with one another ($d_{\text{av,inter}}[\text{Mo}-\text{Mo}] = 3.68$ Å (telluride); 3.31 Å (selenide)). The Ni atoms occupy interstices in the chalcogen atom network and are disordered over pairs of inversion-related sites (occupancy: 0.423(7) and 0.33(-); separation: 2.586(6) Å and 2.279(16) Å for the telluride and selenide,² respectively). Their chalcogen atom environments are shown in Fig. 2 and are compared to those of the ternary metal constituents in other triclinic Chevrel phases.

While the Ba atom in BaMo_6S_8 (10) has eight nearest chalcogen neighbors which form a distorted cube, the Fe atoms in $\text{Fe}_2\text{Mo}_6\text{S}_8$ (11) have five which form a distorted square-pyramid, and the Cu atoms in $\text{Cu}_{1.8}\text{Mo}_6\text{S}_8$ (12) and the Ni atoms in both $\text{Ni}_{0.66}\text{Mo}_6\text{Se}_8$ (8) and $\text{Ni}_{0.84}\text{Mo}_6\text{Te}_8$ each have four which form distorted tetrahedra. Except for Ba all metal atoms are distributed over pairs of inversion-related sites. In the Cu- and Fe-containing compounds these pairs are centered at the cell origin (corresponding to two of the six so-called

“inner sites” of the rhombohedral structure (13)), and they are far away from the Mo_6 cluster (centered at 1/2, 1/2, 1/2). In the Ni compounds these pairs are centered at the midpoint of one of the cell axes (corresponding to two of the six so-called “outer sites” of the rhombohedral structure), and they are close to the Mo_6 cluster. As shown in Fig. 1 the Ni atoms interact strongly with one Mo_6 cluster via two bonds ($d[\text{Ni}-\text{Mo}] = 2.603(3)$ Å and 2.958(3) Å (telluride); 2.420(10) Å and 2.824(10) Å (selenide)), and weakly with another via one bond ($d[\text{Ni}-\text{Mo}] = 2.985(3)$ Å (telluride); 2.862(10) Å (selenide)). Pauling bond order (PBO) sums calculated along the lines described before (14, 15) for the various metal-metal bonds in these compounds are summarized in Table III. As expected the contributions of the molybdenum-metal bonds, ($\text{Mo}-M_x$) to the PBO sums are large only for compounds in which the M atoms occupy outer sites, i.e., for those containing Ni, while they are practically zero for compounds in which the M atoms occupy inner sites, i.e., for $\text{Fe}_2\text{Mo}_6\text{S}_8$ (11) and $\text{Cu}_{1.8}\text{Mo}_6\text{S}_8$ (12), and also for BaMo_6S_8 (10). The difference between the PBO sums of the Mo-Mo bonds of the Ni-based ternary compounds and those of the corresponding Ni-free binary compounds is larger for the selenide (1.6

² Calculated from the lattice constants and atomic parameters reported in Ref. (8).

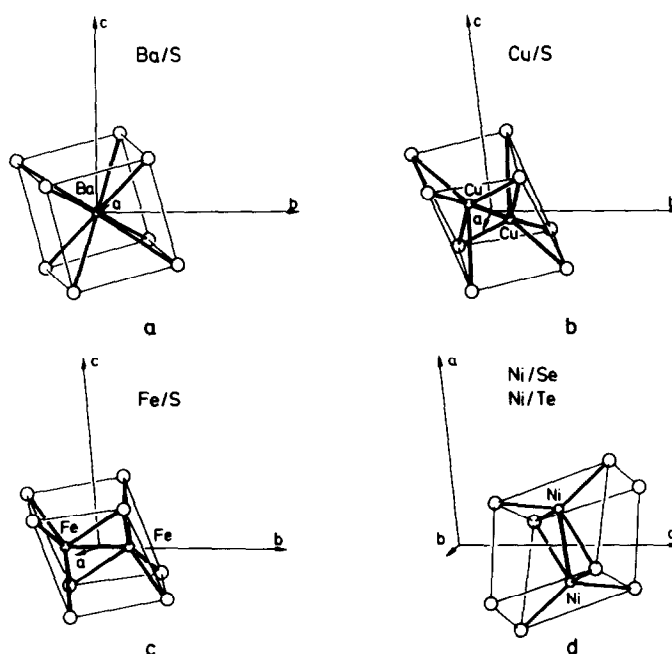


FIG. 2. Comparison between chalcogen atom environments in various triclinic Chevrel phases. (a) BaMo_6S_8 (Ba/S), (b) $\text{Cu}_{1.8}\text{Mo}_6\text{S}_8$ (Cu/S), (c) $\text{Fe}_2\text{Mo}_6\text{S}_8$ (Fe/S), (d) $\text{Ni}_{0.66}\text{Mo}_6\text{Se}_8$ (Ni/Se) and $\text{Ni}_{0.85}\text{Mo}_6\text{Te}_8$ (Ni/Te). In each compound the pseudo-trigonal axis runs along $[111]$.

electron pairs) than for the telluride (0.2 electron pair). This is consistent with a previously observed trend concerning the apparent contraction of the Mo_6 clusters

due to charge transfer (13) and matrix effects (14) as one goes from the sulfides to the selenides and tellurides. The relatively large PBO contribution of the metal-metal

TABLE III
PAULING-BOND-ORDER^a SUMS FOR METAL-METAL BONDS IN TRICLINIC AND RHOMBOHEDRAL CHEVREL PHASES

Compound	$\Sigma(\text{Mo-Mo})^b$	$\Sigma(\text{Mo-M}_x)^c$	$\Sigma(M_x-M_x)$	$\Sigma + \Sigma + \Sigma$	Ref.
Mo_6S_8	14.63	—	—	14.63	(16)
BaMo_6S_8	19.32	0.17	—	19.49	(10)
$\text{Fe}_2\text{Mo}_6\text{S}_8$	19.61	0.14	1.16	20.91	(11)
$\text{Cu}_{1.8}\text{Mo}_6\text{S}_8$	16.58	0.09	0.14	16.81	(12)
Mo_6Se_8	15.11	—	—	13.11	(9)
$\text{Ni}_{0.66}\text{Mo}_6\text{Se}_8$	16.69	0.25	2.20	19.14	(8)
Mo_6Te_8	15.62	—	—	15.62	(4)
$\text{Ni}_{0.85}\text{Mo}_6\text{Te}_8$	15.43	0.12	1.48	17.03	(this work)

^a Calculated per formula unit by using the relation $d(n) = d(1) - 0.60 \log n$ (n = bond order) and the following atomic radii (\AA) for bond order of unity: 1.3095 (Mo), 2.011 (Ba), 1.185 (Fe), 0.955 (Cu), and 1.156 (Ni).

^b Includes 12 Mo-Mo intracluster and 6 Mo-Mo intercluster bonds.

^c Includes Mo-M bonds up to 3.5 \AA .

bonds, M_x-M_x , in the Fe compound compared to that in the Cu compound is presumably due to the choice of the single-bond radii (Fe: 1.185 Å, Cu: 0.955 Å). Pair formation of the metal atoms has so far only been ascertained for those compounds in which these atoms occupy inner sites, i.e., for $\text{Fe}_2\text{Mo}_6\text{S}_8$ ($PP(\text{Fe}) = 1.0$, $d[\text{Fe}-\text{Fe}] = 2.51$ Å (11)), and $\text{Cu}_{1.8}\text{Mo}_6\text{S}_8$ ($PP(\text{Cu}) = 0.90$, $d[\text{Cu}-\text{Cu}] = 2.54$ Å (12)).

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References

1. R. CHEVREL AND M. SERGENT, in "Superconductivity in Ternary Compounds," Vol. 1, Chapter 2 (Ø. Fischer and M. B. Maple, Eds.), Springer, Berlin 1982.
2. W. HÖNLE AND K. YVON, *Z. Kristallogr.* **162**, 103 (1983).
3. W. HÖNLE, H. G. VON SCHNERING, A. LIPKA, AND K. YVON, *J. Less-Common Met.* **71**, 135 (1980).
4. W. HÖNLE AND H. G. VON SCHNERING, in preparation.
5. G. M. SHELDRIK, SHELXTL, CAMBRIDGE, 1976, unpublished.
6. E. PARTHÉ AND L. GELATO, *Acta Crystallogr. Sect. A* **40**, 169 (1984).
7. "Superconducting Quantum Interference Device," S.H.E. Corp., San Diego (1981).
8. O. BARS, J. GUILLEVIC, AND D. GRANDJEAN, *J. Solid State Chem.* **6**, 335 (1973).
9. O. BARS, J. GUILLEVIC, AND D. GRANDJEAN, *J. Solid State Chem.* **6**, 48 (1973).
10. R. BAILLIF, A. DUNAND, J. MULLER, AND K. YVON, *Phys. Rev. Lett.* **47**, 672 (1981).
11. K. YVON, R. CHEVREL, AND M. SERGENT, *Acta Crystallogr. Sect. B* **36**, 685 (1980).
12. K. YVON, R. BAILLIF, AND R. FLÜKIGER, *Acta Crystallogr. Sect. B* **35**, 2859 (1979).
13. K. YVON, in "Current Topics in Materials Science" (E. Kaldis, Ed.), Vol. 3, Chap. 2, North Holland, Amsterdam (1979).
14. J. D. CORBETT, *J. Solid State Chem.* **39**, 56 (1981).
15. W. HÖNLE, H. D. FLACK, AND K. YVON, *J. Solid State Chem.* **49**, 157 (1983).
16. R. CHEVREL, M. SERGENT, AND J. PRIGENT, *Mater. Res. Bull.* **9**, 1487 (1974).